

 $KCI:Ag^{+}$   $\int_{D} B A$   $23^{9K}$  P[001]  $100 \frac{kp}{cm^{2}}$   $50^{9K}$  P[001]  $100 \frac{kp}{cm^{2}}$   $50^{9K}$ 

corption constant K and its change  $\Delta K$  unstress versus photon energy for KCl:Ag<sup>\*</sup>

of aggregate centres the samfor a few minutes at about the melting point and then oom temperature.

part of Fig. 1 shows the UV NaCl:Cu<sup>+</sup> at 77 °K. The main eV, called the D-band, correansition  $\Gamma_1^+ \rightarrow \Gamma_5^+$ . In the small 4.5 eV, called the A-band, the ons  $\Gamma_1^+ \rightarrow \Gamma_3^+$  and  $\Gamma_1^+ \rightarrow \Gamma_5^-$ .  $\Gamma_1^+ \rightarrow \Gamma_3^+$ ,  $\Gamma_5^+$ , contribute to the  $\Gamma_{\pm}^+$  denotes the transformation i the unperturbated electronic defect under the elements of oup O<sub>h</sub>. When uniaxial stress

oscillator strength  $f_{||} - f_{\perp}$  normalized to r strength f plotted versus temperature is applied to the crystal the absorption constants  $K_{\parallel}$  and  $K_{\perp}$  measured with light polarized parallel and perpendicular to the stress axis are different. This difference  $\Delta K = K_{\parallel} - K_{\perp}$  is plotted versus spectral energy for [001] stress in the lower part of Fig. 1. A comparison of  $\Delta K(E)$  with the absorption spectrum K(E) in the upper part of Fig. 1 shows that the change of the oscillator strength under applied stress is the dominant effect, whereas a shift of the bands due to a splitting of the excited states is negligible.  $\Delta K$  is a linear function of the stress. We also note that  $\Delta K$  differs in sign for the two bands.

Fig. 3 summarizes the temperature dependence of the stress effects. The oscillator strength f(T) for a single band is found from a Gauss-curve analysis using the data of Fußgänger [1]. We find that  $\Delta f/f = (f_{\parallel} - f_{\perp})/f$  has the same sign for corresponding bands in NaCl:Cu<sup>+</sup> and KCl:Ag<sup>+</sup>. To explain sign and temperature dependence of  $\Delta f/f$  we will discuss three effects: 1. a frequency splitting of the far IR resonance mode, 2. a small off-centre distortion of the defect, and 3. a quadratic electron-lattice interaction. We will show in the discussion of the experimental data that all three effects contribute to the observed dichroism, but the main contribution comes from the first and second effects mentioned above.

## 3. Theory

To investigate the interaction between the electronic states of the defect and the ionic displacements, it is convenient to classify the electronic wave functions and the normal coordinates Q(l) of the perturbed lattice according to their transformation properties under the elements of the point group of the defect site. For well localized defects it is sufficient to consider the interaction within the octahedral complex  $Ag^+(Cl^-)_6$  or  $Cu^+(Cl^-)_6$  and introduce the symmetry coordinates  $q_{xi}$  of the complex which are basis functions of the irreducible representations  $\Gamma_{\alpha}$  of  $O_h$  (*i* denotes the row index). Small deviations from  $O_h$ -symmetry due to lattice vibrations and static distortions (Jahn-Teller effect, off-centre displacement) of the lattice cell will be treated as perturbations. The symmetry coordinates  $q_{xi}$  are linear combinations of the normal coordinates Q(l) of the perturbed lattice:

$$q_{\alpha i} = \sum_{i} c_{\alpha i}(l) Q(l) . \tag{1}$$

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Each symmetry coordinate is associated with a distortion with a potential energy  $\frac{1}{2} \omega_x^2 q_{xi}^2$  where  $\omega_x$  is an effective frequency of the complex.

The Hamiltonian H of the system

$$H = H_{\rm e} + H_{\rm l} + H_{\rm e\,l} \tag{2}$$

consists of the Hamiltonian of the defect electrons  $H_{\rm e}$ , the lattice Hamiltonian  $H_{\rm l}$ , and the operator of the electron-lattice interaction  $H_{\rm el}$ . We expand the electron lattice interaction in the Hamiltonian in terms of the symmetry coordinates  $q_{\rm xi}$ , taking linear (equation (3a)) and quadratic terms (equation (3b)) into account:

$$H_{c1} = H_1 + H_2 \,, \tag{3}$$

$$H_1 = \sum_{i} V_{xi}(\mathbf{r}) q_{xi}^{-}, \qquad (3a)$$

$$H_2 = \sum_{\alpha, i} \sum_{\beta, j} V_{\alpha i \beta j}^-(\mathbf{r}) q_{\alpha i}^- q_{\beta j}^+ + \cdots .$$
(3b)

r denotes the set of electronic coordinates and the upper sign the parity of the distortion.  $V_{\alpha i}^{-}(r)$  has the same transformation properties as the *i*-th basis func-

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